# FORTRAN IMPLEMENTATION FOR NUCLIDE LIQUID DROP MODEL CALCULATIONS WITH NUMERICAL TECHNIQUES

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ABSTRACT: This research proposes a numerical variation method developed to find the best parameters of the nuclide liquid drop model in calculating its mass. This method was developed because there was no systematic method for finding empirical model parameters in calculating the nuclide mass or binding energy. The modelling used in this study uses the Fortran programming language or formula translation. The results obtained provide a very significant improvement in the model for calculating the mass of stable nuclides. The closeness of the empirical model calculation results to the experimental results shows the model's validity. The delta deviation results from the proposed numerical method give the smallest value compared to existing methods at 111. Meanwhile, the error rate in the proposed method is 0.00088% or 8.84 x 10-6.

**KEYWORDS**: Numerical; Nuclide; Fortran; Proton; Neutrons

## 1.0 INTRODUCTION

The liquid drop model for calculating nuclide binding energies has been developed since 1913 and was first proposed by Bohr [1], [2]. The Bohr model uses many assumptions and approximations to simplify the description of nuclide behaviour. One example is the Bohr model of the hydrogen atom, which only considers one electron, while nuclides have many protons and neutrons that interact in complex

ways [3], [4]. The binding energy formulated on a semi-empirical basis by Weizsacker turns out to be sufficient to provide an estimate for knowing the energy balance of various nuclear quantities such as nuclide mass, binding energy, energetic beta decay, alpha decay, nuclear reactions, resonant reaction cross sections and energetic fission [1], [5]. This model has undergone improvements with various choices of model parameters. The empirical model used assumes that the binding energy consists of nucleon binding energy, which is proportional to the mass number A, surface binding energy, considering that liquid drops always have surface energy, Coulomb repulsion energy, which is directly proportional to the atomic number Z squared, symmetry energy to the number of neutrons and protons, and the discrete energy due to nucleons as fermion particles. [1], [5].

This research aims to create a numerical model to calculate nuclide energy so minimal errors occur. [6]. The method assumes that each parameter is referred to sequentially as a1, a2, a3, a4, and a5. Hence, the key is to set the parameters so that the error in the model calculation results compared to experimental data is the smallest. The determination is carried out using a trial system, and it is unclear what the best trial method is to obtain the best parameter values. [7], [8], [9], [10]. The modelling used in this study uses the Fortran programming language or formula translation. The results of the numerical calculations are compared with the parameters that have been tested, including those given by Segre. [11], [12], Elton [13], Kaplan[14], Godfrey [15], Vincent [16], and Keenan [17]. This paper proposes an iterative numerical variation method to obtain the best parameter quantities. [9], [16]. The selection of nuclide data is limited to stable nuclides only.

## 2.0 METHOD

Nuclides with atomic number Z have Z protons. [18]. The mass number A indicates the presence of A nucleons (protons and neutrons)—the number of nucleons that comprise the nuclide as a constituent with Equation 1.

$$M = Z.M_P + (A - Z).M_N = 1.008983 A - 0.000839 Z....(1)$$

The mass of a proton is 1.0081437, and the mass of a neutron is 1.0089830 mass units. However, the mass changes according to the

binding energy used because of the negative binding energy. Using the water drop model, the first correction comes from the binding energy of nucleons (protons or neutrons) whose magnitude is proportional to A and is given by Equation 2.

$$M_V = a_V A \dots (2)$$

The constant a1 is a constant that needs to be determined, as are the constants involved in subsequent corrections. This first correction ignores the situation where, at the surface, the nucleon binding energy is smaller. Therefore, it is necessary to provide a surface effect correction, the amount of which is calculated using Equation 3.

$$M_S = a_{S.} A^{2/3}$$
....(3)

The next correction comes from the mutual repulsion force between protons, which is proportional to Z(Z-1). Based on quantum calculations, this energy is only proportional to  $Z_2$ , so the correction can be calculated with equation 4.

$$M_C = a_{C.}Z(Z-1)/A^{1/3}$$
....(4)

For a charge with a uniform distribution, the constant  $a_3$  is 3e2/5.r0 = 0.000763. The last fact shows stability around the position where the number of protons and neutrons equals. This correction is given by equation 5.

$$M_{SY} = a_{SY}(A/2-Z)^2/A$$
 .....(5)

All calculated terms show a continuous dependence of the nuclide energy on the number of neutrons and protons. However, some components are still not continuous, namely due to the evenness of A. This energy correction is given by equation 6.

$$D(A, Z) = + a\delta f(A)$$
....(6)

for A even Z odd, = 0 for A odd and = -  $a\delta$  f(A) for A even Z odd.

Various experts, among others, have established the approach for unknown coefficients, which is given by reference. Next, the form of Equation 7 will be used.

$$M = a_1 A + a_2 Z + a_3 A_2/3 + a_4 (A - 2Z) + a_5 Z(Z - 1)_2 + a_6 f(A)_1 (7)$$

Where:

f(A) = 1/A for references 1 and 2.

f(A) = 1/A3/4 for references 3, 4, 5, 6.

f(A) = 1/A1/2 for reference 7.

The Least Numerical Variation Method is a computational physics and chemistry technique used to estimate the basic energy (ground state energy) of quantum systems such as atoms, molecules or nuclides. This method is generally used when analytical solutions are unavailable or impractical. This method is based on the variational principle, where the system energy is estimated by varying the trial wave function used in the calculations. The main goal is to find the test wave function that gives the lowest energy (basic energy).

Numerical approaches are generally used to find the smallest variations. This method involves numerical techniques such as numerical integration, solving partial differential equations, and numerical optimisation to calculate the energy expectation value in the context of quantum mechanics. The advantage of this numerical method is that it can be used for complex quantum systems where analytical calculations are impossible. In addition, this method can be adapted to various situations, including using different basis sets to improve the accuracy of the results.

This method varies the value of i from 1 to 6. The variations start with  $a_1$ . The initial value taken is the price given by Kaplan. [19].  $a_1$  is varied numerically to obtain the lowest value of the closeness of the experimental results to the calculation results of Equation 8.

$$delta = \sum_{i=1} | M_{i(exp)} - M_{i(model)} | \dots$$
 (8)

The variation of a1 is driven from the initial value  $a_1\omega$ , whose magnitude is a1-nh, where h is the step length, set h = 0.00001 and n taken 1000 times. The lowest delta price is taken by moving  $a_1\omega = a_1\omega + h$  by 2n steps. After the best a1 price is obtained, the same procedure is carried out to find other ai prices. After all the best ai are obtained, the procedure can be repeated using a smaller h value for the second iteration.

Nuclide stability tests can also be carried out based on the water drop model by looking for the relationship between Z based on the requirement that the differential of the binding energy concerning Z (with A constant) is equal to zero. This will give the second equation given by equation 9.

$$Z = A / \{1.98 + 0.015 A2/3\}$$
 .....(9)

For numerical purposes, it is sufficient to carry out numerical variations of Z (fixed A) and take the value of Z that gives the lowest binding energy.

## 3.0 RESULT AND DISCUSSION

# 3.1 Delta Comparison of Liquid Drop Model

In Table 1, a comparison of delta is shown as the cumulative deviation price against the experimental price. From the results of this comparison, it can be seen that the numerical variation method provides the smallest deviation with parameter values that are better than the parameters given by several researchers. Meanwhile, the relationship between Z and A for stable nuclides, as required by Segre, is closer to the truth for mass numbers below 20. In this study, the Least Numerical Variation Method gives the smallest delta value of 1.0451118, indicating a powerful approach for calculating the fundamental energy of quantum systems. These results are close to the values predicted from standard theory for similar systems, indicating satisfactory accuracy in modelling the quantum properties of the generated energy. The use of the Least Numerical Variation Method in

this research is justified by its ability to adapt to system complexity and flexibility in selecting the appropriate test wave function [20], [21], [22].

Table 1. Comparison of Liquid Drop Model Coefficients

|   |             |             |             |             | 1           |             |             |  |
|---|-------------|-------------|-------------|-------------|-------------|-------------|-------------|--|
| а | Budi *      | Segre       | Elton       | Kaplan      | Godfrey     | Vincent     | Keenan      |  |
| 1 | 0.9935337   | 0.9917550   | 0.9917400   | 0.9939100   | 0.9917400   | 0.9918700   | 0.9938400   |  |
| 2 | - 0.0008502 | 0.0008400   | - 0.0008390 | - 0.0008500 | - 0.0008390 | - 0.0008390 | - 0.0008390 |  |
| 3 | 0.0141000   | 0.0191100   | 0.0141000   | 0.0140000   | 0.0191150   | 0.0191270   | 0.0139614   |  |
| 4 | 0.0209428   | 0.1017500   | 0.0210000   | 0.8300000   | 0.1017500   | 0.0254526   | 0.0204051   |  |
| 5 | 0.0006300   | 0.0007630   | 0.0006300   | 0.0006270   | 0.0007626   | 0.0007625   | 0.0006390   |  |
| 6 | - 0.0150551 | - 0.0130000 | - 0.1450000 | - 0.0360000 | - 0.0360840 | - 0.0365140 | - 0.0359774 |  |
|   |             |             |             |             |             |             |             |  |
|   |             |             |             |             |             |             |             |  |
|   |             |             |             |             |             |             |             |  |
| n |             |             |             |             |             |             |             |  |
| Δ | 1.0451118   | 67.0478400  | 10.4621600  | 49.8922900  | 49.8064600  | 2.1128820   | 7.7187860   |  |

## 3.2 Comparison with Other Methods

Table 2 presents an example of a comparison between the various parameters used. The numerical variation method provides calculation results closer to the experimental results. The small deviation or deviation from the experimental values leads to the conclusion that the liquid drop model for calculating the mass of nuclides is quite representative. Thus, estimates of the energy released from a nuclear reaction can be made if the nuclides resulting from the response are known. From Table 2, it can be seen that the proposed research results produce values close to the experimental calculation results. The error value given by the method proposed in this research is 0.00088% or 8.84 x10-6.

Nuclide droplets, or specific nuclides, have important applications in medicine, industry, energy, research, and the environment. In medicine, radioactive isotopes such as technetium-99m are used for diagnostic imaging, while iodine-131 aids in thyroid cancer therapy. In industry, nuclides are used for leak detection, radiography, and measuring the thickness of materials. In energy, isotopes such as uranium-235 and plutonium-239 are used as fuel for nuclear reactors, while plutonium-238 is used in radioisotope thermoelectric generators for space missions. Nuclides also play a role in research, such as

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radiometric dating with carbon-14 and tracking biological reactions. In the environment, isotopes monitor pollution and track water flows. Additionally, nuclides support security through explosives detection and cargo screening. These broad applications demonstrate the importance of nuclides in modern technology.

Table 2. Comparison between various calculations and experiments

|    |     |         |         |         | i various calculations and experimen |         |         |         |         |
|----|-----|---------|---------|---------|--------------------------------------|---------|---------|---------|---------|
| Z  | A   | EXP     | PRESENT | SEGRE   | ELTON                                | KAPLAN  | GODFREY | VINCENT | KEENAN  |
| 68 | 164 | 164.930 | 164.930 | 165.479 | 165.000                              | 165.348 | 165.363 | 164.941 | 164.981 |
| 67 | 165 | 165.930 | 165.931 | 166.451 | 166.001                              | 166.325 | 166.334 | 165.942 | 165.984 |
| 68 | 166 | 166.932 | 166.932 | 167.477 | 167.002                              | 167.346 | 167.360 | 166.943 | 166.984 |
| 68 | 167 | 167.932 | 167.933 | 168.505 | 168.005                              | 168.370 | 168.388 | 167.945 | 167.986 |
| 68 | 168 | 167.934 | 167.933 | 168.399 | 168.004                              | 168.281 | 168.279 | 167.945 | 167.987 |
| 70 | 168 | 168.934 | 168.933 | 169.475 | 169.005                              | 169.344 | 169.356 | 168.945 | 168.986 |
| 69 | 169 | 169.936 | 169.936 | 170.562 | 170.009                              | 170.418 | 170.446 | 169.949 | 169.990 |
| 68 | 170 | 169.935 | 169.934 | 170.449 | 170.007                              | 170.322 | 170.328 | 169.947 | 169.989 |
| 70 | 170 | 170.936 | 170.935 | 171.473 | 171.007                              | 171.342 | 171.353 | 170.947 | 170.989 |
| 70 | 171 | 171.936 | 171.936 | 172.501 | 172.010                              | 172.366 | 172.381 | 171.949 | 171.991 |
| 70 | 172 | 172.938 | 172.938 | 173.528 | 173.011                              | 173.388 | 173.407 | 172.950 | 172.991 |
| 70 | 173 | 173.939 | 173.939 | 174.557 | 174.013                              | 174.412 | 174.437 | 173.953 | 173.994 |
| 70 | 174 | 173.940 | 173.938 | 174.447 | 174.013                              | 174.320 | 174.324 | 173.952 | 173.995 |
| 72 | 174 | 174.941 | 174.939 | 175.526 | 175.013                              | 175.386 | 175.404 | 174.952 | 174.994 |
| 71 | 175 | 175.943 | 175.943 | 176.616 | 176.018                              | 176.461 | 176.495 | 175.957 | 175.998 |
| 70 | 176 | 175.942 | 175.940 | 176.499 | 176.015                              | 176.363 | 176.375 | 175.954 | 175.997 |
| 72 | 176 | 176.943 | 176.941 | 177.524 | 177.016                              | 177.384 | 177.401 | 176.954 | 176.997 |
| 72 | 177 | 177.944 | 177.943 | 178.553 | 178.019                              | 178.408 | 178.429 | 177.957 | 177.999 |
| 72 | 178 | 178.946 | 178.944 | 179.580 | 179.020                              | 179.430 | 179.456 | 178.958 | 179.000 |
| 72 | 179 | 179.947 | 179.946 | 180.610 | 180.022                              | 180.455 | 180.486 | 179.961 | 180.002 |
| 72 | 180 | 179.947 | 179.945 | 180.497 | 180.021                              | 180.360 | 180.370 | 179.960 | 180.003 |
| 74 | 180 | 180.948 | 180.946 | 181.578 | 181.022                              | 181.428 | 181.453 | 180.960 | 181.002 |
| 73 | 181 | 181.948 | 181.947 | 182.550 | 182.024                              | 182.404 | 182.423 | 181.962 | 182.005 |
| 74 | 182 | 182.950 | 182.948 | 183.576 | 183.025                              | 183.426 | 183.449 | 182.963 | 183.005 |
| 74 | 183 | 183.951 | 183.949 | 184.606 | 184.028                              | 184.451 | 184.479 | 183.965 | 184.008 |
| 74 | 184 | 183.953 | 183.950 | 184.497 | 184.028                              | 184.359 | 184.366 | 183.966 | 184.010 |
| 76 | 184 | 184.953 | 184.950 | 185.575 | 185.028                              | 185.424 | 185.446 | 184.965 | 185.008 |
| 75 | 185 | 185.954 | 185.953 | 186.664 | 186.032                              | 186.499 | 186.537 | 185.969 | 186.011 |
| 74 | 186 | 186.956 | 186.952 | 187.574 | 187.031                              | 187.423 | 187.443 | 186.968 | 187.011 |
| 76 | 187 | 187.956 | 187.954 | 188.602 | 188.034                              | 188.447 | 188.472 | 187.971 | 188.014 |
| 76 | 188 | 188.958 | 188.955 | 189.630 | 189.035                              | 189.469 | 189.499 | 188.971 | 189.014 |
| 76 | 189 | 189.959 | 189.957 | 190.660 | 190.038                              | 190.495 | 190.529 | 189.974 | 190.017 |
| 76 | 190 | 190.961 | 190.957 | 191.628 | 191.038                              | 191.468 | 191.496 | 190.974 | 191.017 |
| 77 | 191 | 191.961 | 191.960 | 192.720 | 192.042                              | 192.544 | 192.589 | 191.978 | 192.021 |
| 76 | 192 | 192.963 | 192.960 | 193.686 | 193.042                              | 193.516 | 193.554 | 192.978 | 193.021 |
| 77 | 193 | 193.963 | 193.961 | 194.656 | 194.044                              | 194.491 | 194.522 | 193.980 | 194.023 |
| 78 | 194 | 194.965 | 194.963 | 195.684 | 195.045                              | 195.514 | 195.550 | 194.981 | 195.024 |
| 78 | 195 | 195.965 | 195.964 | 196.715 | 196.048                              | 196.539 | 196.581 | 195.984 | 196.027 |
| 78 | 196 | 195.966 | 195.964 | 196.600 | 196.048                              | 196.443 | 196.462 | 195.984 | 196.028 |
| 80 | 196 | 196.967 | 196.965 | 197.683 | 197.049                              | 197.512 | 197.547 | 196.984 | 197.027 |
| 79 | 197 | 197.968 | 197.968 | 198.776 | 198.053                              | 198.590 | 198.642 | 197.988 | 198.030 |
| 78 | 198 | 197.967 | 197.967 | 198.654 | 198.051                              | 198.488 | 198.517 | 197.986 | 198.030 |
| 80 | 198 | 198.968 | 198.968 | 199.681 | 199.052                              | 199.510 | 188.544 | 198.987 | 199.031 |
| 80 | 199 | 199.968 | 199.969 | 200.711 | 200.055                              | 200.535 | 200.574 | 199.990 | 200.033 |
| 80 | 200 | 200.970 | 200.971 | 201.740 | 201.056                              | 201.559 | 201.602 | 200.991 | 201.034 |
| 80 | 201 | 201.971 | 201.973 | 202.771 | 202.059                              | 202.585 | 202.633 | 201.994 | 202.037 |
| 80 | 202 | 202.972 | 202.974 | 203.738 | 203.059                              | 203.557 | 203.599 | 202.994 | 203.037 |
| 81 | 203 | 203.974 | 203.977 | 204.833 | 204.064                              | 204.637 | 204.696 | 203.998 | 204.041 |
| 80 | 204 | 203.973 | 203.975 | 204.709 | 204.062                              | 204.533 | 204.568 | 203.997 | 204.040 |
| 82 | 204 | 204.974 | 204.977 | 205.798 | 205.064                              | 205.607 | 205.659 | 204.998 | 205.041 |
| 81 | 205 | 205.974 | 205.978 | 206.767 | 206.066                              | 206.581 | 206.626 | 206.000 | 206.044 |
| 82 | 206 | 206.976 | 206.980 | 207.796 | 207.067                              | 207.605 | 207.655 | 207.001 | 207.044 |
| 82 | 207 | 207.977 | 207.982 | 208.828 | 208.070                              | 208.632 | 208.687 | 208.004 | 208.047 |
|    | 207 | -0.077  | 20,.702 | 200.020 | 200.070                              | 200.002 | 200.007 | 200.001 | 200.047 |

## 4.0 CONCLUSION

The small deviation or deviation from the experimental values suggests that the liquid drop model for calculating the mass of nuclides is quite representative. Thus, estimates of the energy released from a nuclear reaction can be made if the nuclides resulting from the response are known. The magnitude of the error delta in the numerical method proposed in this study is 1.00. This value gives an error rate for research results of 0.00088%, equivalent to 8.84x10-6. Suggestions for future research can develop this numerical method for other applications.

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